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Investigation of Turbulent Recovery Factor in Hypersonic Helium Flow

DAVID H. RUDY* AND LEONARD M. WEINSTEIN*
NASA Langley Research Center, Hampton, Va.

Nomenclature

M = Mach number
 r = recovery factor $(T_{aw} - T_e)/(T_o - T_e)$
 R = Reynolds number
 T = temperature

Subscripts

e = local conditions at edge of boundary layer
 aw = adiabatic wall
 o = total
 x = distance along plate from leading edge
 v = virtual origin (at peak of recovery factor in transition region)

AN accurate knowledge of the temperature recovery factor is important in the reduction of wind-tunnel heat-transfer measurements to coefficient form, especially when those measurements are made at nearly adiabatic conditions. Numerous experimental measurements of the re-

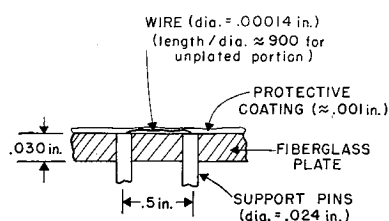


Fig. 1 Lateral cut showing instrumentation for surface temperature measurement (not to scale).

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* Aerospace Engineers, Flow Analysis Section, Aero-Physics Division.

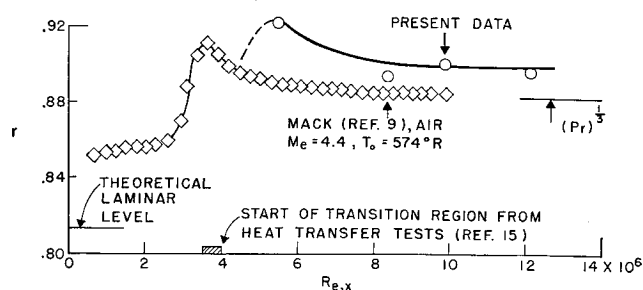


Fig. 2 Recovery factor distribution for wedge model ($M_e = 6.8$, helium).

covery factor have been made for laminar, transitional, and turbulent boundary layers in air for Mach numbers up to 5 (i.e., Refs. 1-9). A few measurements have also been made between Mach 5 and 6.¹⁰⁻¹² There has actually been no pressing need for precise air recovery factor measurements at Mach numbers above 5 because the total temperatures necessary to avoid liquefaction of air at these high Mach numbers place the usual wall temperatures far below the adiabatic values. However, in helium wind tunnels, hypersonic Mach numbers can be achieved without heating the flow and therefore an accurate value of the recovery factor is again needed for reduction of heat-transfer data. The only recovery factor study previously conducted in helium involved laminar and transitional flow at Mach 6.5.¹³ The present investigation was therefore initiated to experimentally obtain the turbulent recovery factor in hypersonic helium flow.

The present measurements of turbulent recovery factor were made in the Langley 22-in. Mach 20 helium tunnel using a thin fiber glass plate mounted on a sharp leading edge (0.002-in. thick), 10° half-angle wedge. The local Mach number M_e on the 10° wedge was 6.8 with a negligible local pressure gradient.

The model was designed so that the fiber glass surface would be essentially insulated and thus, the equilibrium surface temperature would be very nearly the adiabatic-wall temperature. Longitudinal heat conduction was negligible because of the small cross section and low conductivity of the fiber glass plate. Calculations also indicated negligible error as the result of radiation from the tunnel wall. To accurately measure the surface temperature, the model was instrumented with thin wire (0.00014-in. diam) resistance thermometers, the ends of which were copper-plated for attachment to support pins, as shown in Fig. 1. A conservative analysis indicated that conduction through the steel support pins had only a slight (less than 0.1°R) effect on the wire temperature. Therefore, the wires were assumed to be correctly measuring the surface temperature. During the last half of the 93-sec

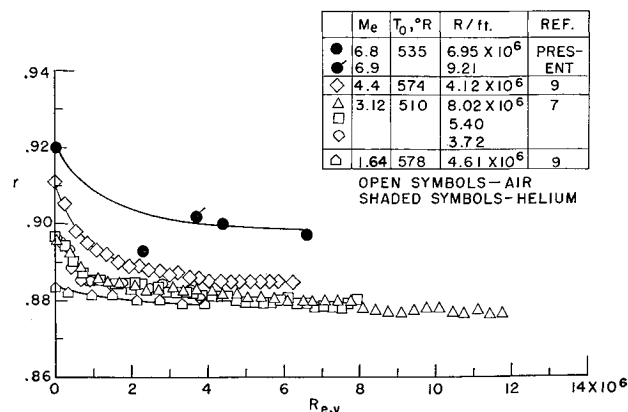


Fig. 3 Variation of turbulent recovery factor with Reynolds number based on distance from recovery factor peak in transition region.

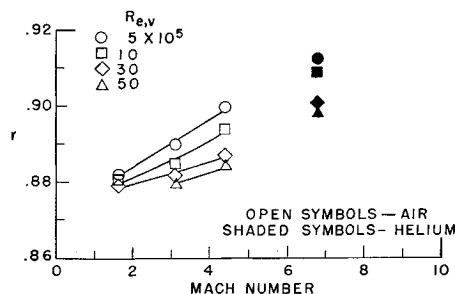


Fig. 4 Effect of Mach number on turbulent recovery factor.

test the ratio of the measured surface temperature to the total temperature changed less than 0.35% for each wire. Recovery factors were then calculated using the equilibrium surface temperature as adiabatic-wall temperature in the defining equation for recovery factor.

The results of the test are shown in Fig. 2. No data were obtained in the laminar region; however, the theoretical values for the helium recovery factor predicted by laminar similar solutions¹⁴ at Mach 6.8 is indicated. The approximate beginning of the transition region determined from heat-transfer tests with unheated flow¹⁵ using a thin metal plate mounted on the same wedge model at similar test conditions is also shown. In Ref. 15 the end of the transition region was found to be located at a distance approximately twice as far from the leading-edge as the beginning. Based on these Ref. 15 results, the first data point is believed to be in the transition region near the location of the peak in recovery factor. This trend is consistent with previous air measurements such as the data of Mack⁹ shown in Fig. 2 for comparison. Based on the fairing through the data, the turbulent helium recovery factor at high Reynolds numbers and Mach 6.8 is approximately 0.899 ± 0.002 . This is significantly higher than the generally used value of 0.883 (cube root of the molecular Prandtl number).

To investigate the effect of Reynolds number and Mach number on the turbulent recovery factor previous air measurements were examined. Fig. 3 shows, for several Mach numbers, the measured recovery factor as a function of Reynolds number based on distance from the location of the peak values in the transition region. As shown in Fig. 3, all the data for each Mach number follow the same trend, with the recovery factor dropping rapidly from a peak value to a slightly decreasing value at higher Reynolds numbers.

The data at several Reynolds numbers are cross plotted in Fig. 4 as a function of Mach number. In contrast to the theory of Tucker and Maslin,¹⁶ the data indicate that the turbulent recovery factor increases with increasing Mach number.

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Nonlinear Elastic Response of an Incased Tubular Grain

ROGER J. EVANS*

AND

MICHAEL S. HAYCOCK†

University of Washington, Seattle, Wash.

Introduction

THE mechanical response of highly loaded solid propellants is significantly nonlinear at strain levels of 5 or 10% in uniaxial stress states. A previous paper in this Journal¹ has shown that useful results may be obtained from the stress analysis of elastic solids that are physically nonlinear conditions of geometric linearity. Solutions were presented for the internal pressurization of a rigidly incased tubular grain and for thermal stresses in an incased grain.

Baltrukonis and Vaishnev² considered finite axisymmetric deformations with regard to the analysis of solid propellant grains. Their work, however, was restricted to a particular class of incompressible elastic solids.

The purpose of this Note is to extend the analysis of Pister and Evans¹ to the more realistic case: where the problem formulation takes into account both physical and geometric nonlinearity. Such an extension is of importance with regard both to extending the range of useful solutions and to justifying the simultaneous consideration of geometric linearity and physical nonlinearity. Solutions for geometric

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* Assistant Professor, Department of Civil Engineering.

† Graduate Student, Department of Civil Engineering.